## Lecture 26: Periodic Motion

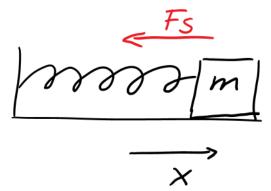
- Motion of a mass at the end of a spring
- Differential equation for simple harmonic oscillation
- Amplitude, period, frequency and angular frequency
- Energetics
- Simple pendulum
- Physical pendulum

#### Mass at the end of a spring

Mass m connected to a spring with spring constant k on a frictionless surface

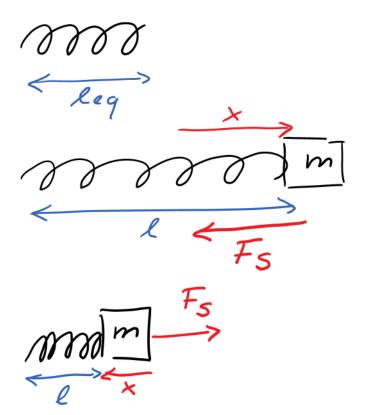
$$F_{x} = -kx$$

Linear restoring spring force



#### **Spring force**

Spring force:  $F_{Sx} = -kx$ 



$$x = l - l_{eq}$$
  
stretch or compression  $k$  force constant

 $F_x$  is negative if x is positive (stretched spring)

 $F_x$  is positive if x is negative (compressed spring)

#### Differential equation of a SHO

Newton's 2<sup>nd</sup> Law:  $\sum F_{\chi} = ma_{\chi}$ 

$$-kx = m\frac{d^2x}{dt^2}$$

$$-\frac{k}{m}x = \frac{d^2x}{dt^2}$$

$$\frac{d^2x}{dt^2} = -\omega^2x$$

$$\omega = \sqrt{\frac{k}{m}}$$

Angular frequency

Differential equation of a Simple Harmonic Oscillator

\*We can always write it like this because m and k are positive

#### **Solution**

$$\frac{d^2x}{dt^2} = -\omega^2 x$$

**Equation for SHO** 

General solution:

$$x = A\cos(\omega t + \varphi)$$

Take 2<sup>nd</sup> derivative: 
$$\frac{d^2x}{dt^2} = -A \omega^2 \cos(\omega t + \varphi) = -\omega^2 x$$

A and  $\varphi$ : two "constants of integration" from solution of a second-order differential equation.

Determined by the initial conditions.

### **Amplitude**

$$x = A\cos(\omega t + \varphi)$$

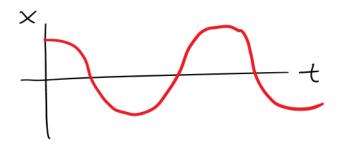
Range of cosine function: -1...+1 
$$\Rightarrow -A \le x(t) \le +A$$

A = Amplitude of the oscillation

#### **Phase Constant**

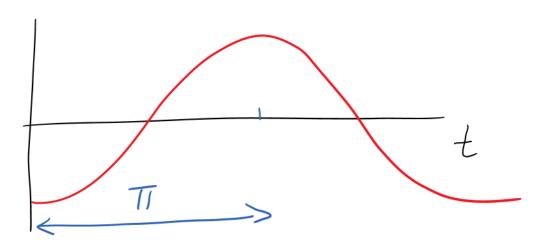
$$x = A\cos(\omega t + \varphi)$$

If 
$$\varphi$$
=0:  
 $x = A \cos(\omega t)$   
 $x(t = 0) = x_0 = A$ 



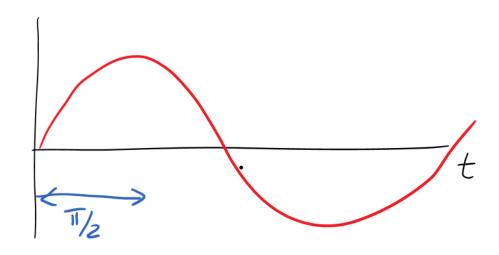
To describe motion with different starting points: Add phase constant to shift the cosine function

$$x = A\cos(\omega t + \varphi)$$



$$X_o = 0$$
:

Shift by  $\frac{\pi}{2}$ 



#### **Initial conditions**

$$x_0 = x(t = 0)$$

$$v_{x0} = v_x(t = 0)$$

$$x_0 = A\cos(0 + \varphi) = A\cos(\varphi)$$
  
$$v_{x0} = -A\omega\sin(0 + \varphi) = -A\omega\sin(\varphi)$$

 $\rightarrow$  two equations for A and  $\varphi$ 

## Position and velocity

$$x = A\cos(\omega t + \varphi)$$

$$v_{\chi} = \frac{dx}{dt} = -A\omega \sin(\omega t + \varphi)$$

At time 
$$t_m$$
:  $x = x_{max} = A$   $\cos(\omega t_m + \varphi) = 1$   $(\omega t_m + \varphi) = 0$  or  $\pi$   $\sin(\omega t_m + \varphi) = 0 \implies v_x(t_m) = 0$ 

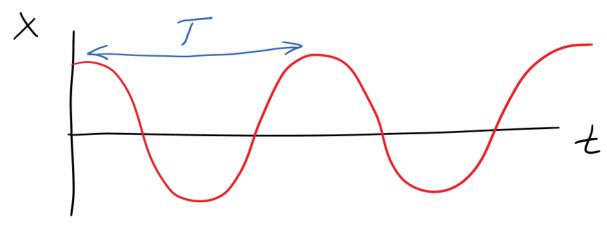
Mass stops and reverses direction when it reaches maximum displacement (turning point)

# **Simulation**

Walter Fendt Spring Pendulum Simulation

## Period and angular frequency

$$x = A\cos(\omega t + \varphi)$$



Time T for one complete cycle: period

 $(\omega t + \varphi)$  changes by  $2\pi$  in time T

$$\omega T = 2\pi \implies \omega = \frac{2\pi}{T} = 2\pi f$$

### Effect of mass and amplitude on period

$$\omega T = 2\pi \implies T = \frac{2\pi}{\omega}$$

$$\omega = \sqrt{\frac{k}{m}} \qquad \Longrightarrow T = \frac{2\pi}{\sqrt{\frac{k}{m}}} \qquad T = 2\pi\sqrt{\frac{m}{k}}$$

$$T = 2\pi \sqrt{\frac{m}{k}}$$

Amplitude A does not appear – no effect on period

Demo: Vertical springs showing effect of *m* and *A* 

### **Energy in SHO**

Potential energy of spring force:  $U = \frac{1}{2}kx^2$ 

at 
$$x=\pm A$$
:
$$U = \frac{1}{2}kA^{2}$$

$$K = 0$$

$$E = \frac{1}{2}kA^{2}$$

$$X = -A$$

$$X = +A$$

$$At x = 0$$
:
$$U = U$$

$$K = K max = E$$

#### **Example**

A block of mass *M* is attached to a spring and executes simple harmonic motion of amplitude *A*. At what displacement(s) *x* from equilibrium does its kinetic energy equal twice its potential energy?

#### **General SHO**

$$\frac{d^2x}{dt^2} = -\omega^2 x$$

**Equation for SHO** 

General solution:

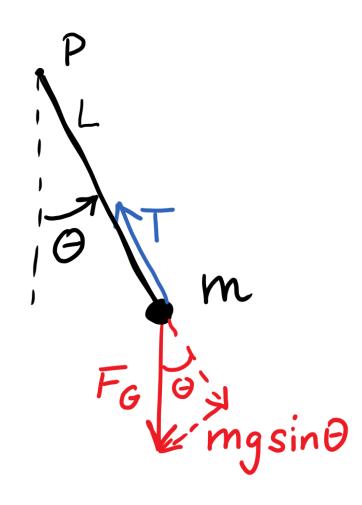
$$x = A\cos(\omega t + \varphi)$$

$$T = \frac{2\pi}{\omega}$$

### Simple Pendulum

Point mass m at the end of a massless string of length L

 $\theta$  = displacement coordinate (with sign) from vertical equilibrium position



#### Simple Pendulum Oscillation

$$\Sigma \tau_z = I \alpha_z$$
 
$$-mg \ L \sin \theta = mL^2 \frac{d^2 \theta}{dt^2}$$

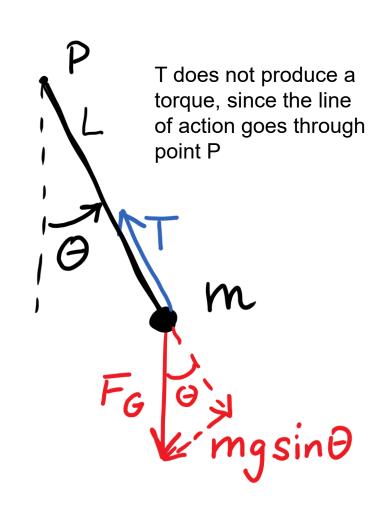
Very complicated differential equation! But for small oscillations:

$$sin\theta \approx \theta$$

And

$$-\frac{g}{L}\theta = \frac{d^2\theta}{dt^2}$$

Differential equation of SHO



### Simple Pendulum Period

$$-\frac{g}{L}\theta = \frac{d^2\theta}{dt^2}$$

Differential equation of simple harmonic oscillator

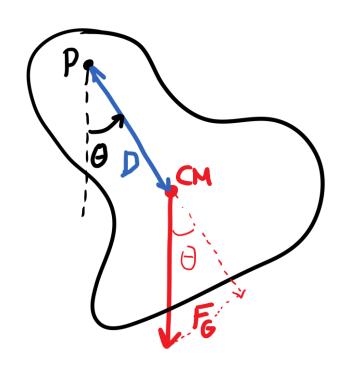
$$\theta(t) = \theta_{max} \cos(\omega t + \varphi)$$

With 
$$\omega = \sqrt{\frac{g}{L}}$$
 and  $T = 2\pi \sqrt{\frac{L}{g}}$ 

$$T = 2\pi \sqrt{\frac{L}{g}}$$

- Period independent of mass
- Period independent of amplitude

## **Physical Pendulum**



Extended object of mass m that swings back and forth about an axis P that does not go through its center of mass CM.

$$\Sigma \tau_z = I \alpha_z$$

$$-mg \ D \ sin\theta = I \frac{d^2 \theta}{dt^2}$$

For small oscillations:  $sin\theta \approx \theta$ 

$$-\frac{mgD}{I}\theta = \frac{d^2\theta}{dt^2}$$

### **Physical Pendulum Oscillation**

$$-\frac{mgD}{I}\theta = \frac{d^2\theta}{dt^2}$$

$$-\frac{mgD}{I}\theta = \frac{d^2\theta}{dt^2} \qquad \text{SHO:} \\ \theta(t) = \theta_{max}\cos(\omega t + \varphi) \\ \omega = \sqrt{\frac{mgD}{I}}$$

$$T = 2\pi \sqrt{\frac{I}{mgD}}$$

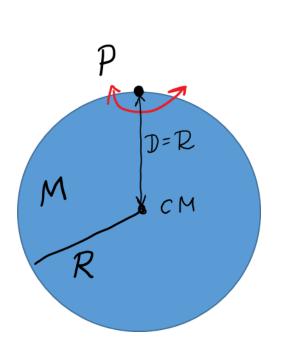
I is moment of inertia about axis P Parallel axis theorem:

$$I_P = I_{CM} + mD^2$$

Demo: Meter stick pivoted at different positions

## **Physicsl Pendulum Example**

A uniform disk of mass M and radius R is pivoted at a point at the rim. Find the period for small oscillations.



$$T = 2\pi \sqrt{\frac{I}{mgD}} \qquad I_P = I_{CM} + mD^2$$